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SPECIAL ISSUE ON

Beating Classical and Quantum Limits in Optics

In order to submit the seminars and lectures in the up-to-date subjects for Iraqi researchers and students, this issue of the Iraqi Journal of Applied Physics (IJAP) is presenting seminar titled "*Beating Classical and Quantum Limits in Optics*" authored by Mankei Tsang.

INVITATION TO PARTICIPATE

To all they would like to submit seminars or scientific lectures during the third semester of the **I.S.A.R.E.S.T.** (July, August and September) in 2007, you are kindly requested to contact the secretary of the **I.S.A.R.E.S.T.** for date and presentation arrangements of the seminars or lectures. Please, do not hesitate to participate in our activities, this chance might be required by young scientists in our country, IRAQ, to develop and grow as well as introduce the professors and experts in field. You could find us on the post address, emails and mobile below:

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Beating Classical and Quantum Limits in Optics

Mankei Tsang

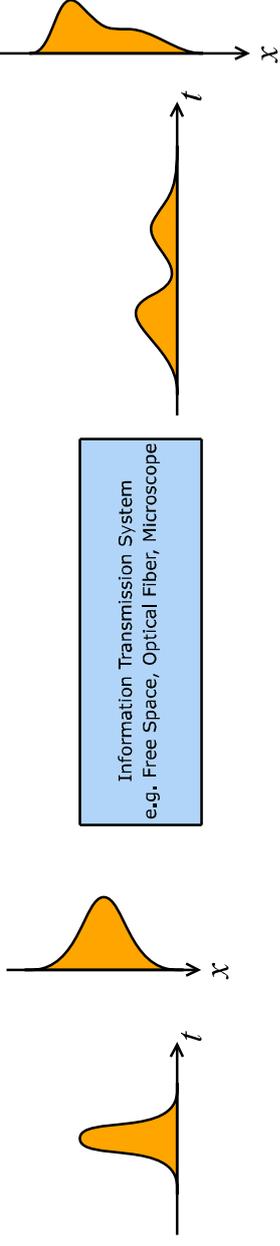
mankei@optics.caltech.edu

Department of Electrical Engineering, Caltech



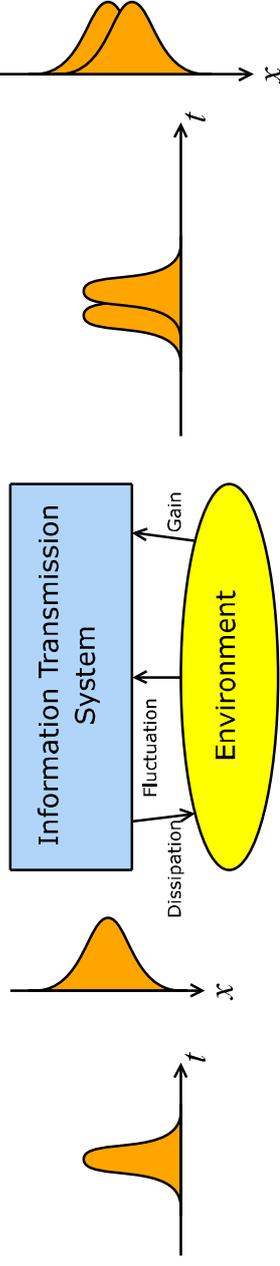
Classical and Quantum Distortions

Classical Effects



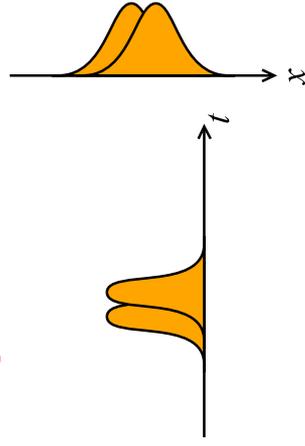
Dispersion, Diffraction, Rayleigh Criterion, Loss, Nonlinearity, ...

Quantum Decoherence



Langevin Noise, Amplified Spontaneous Emission, Gordon-Haus Timing Jitter, ...

Quantum Limits



Standard Quantum Limits, Heisenberg Uncertainty Principle

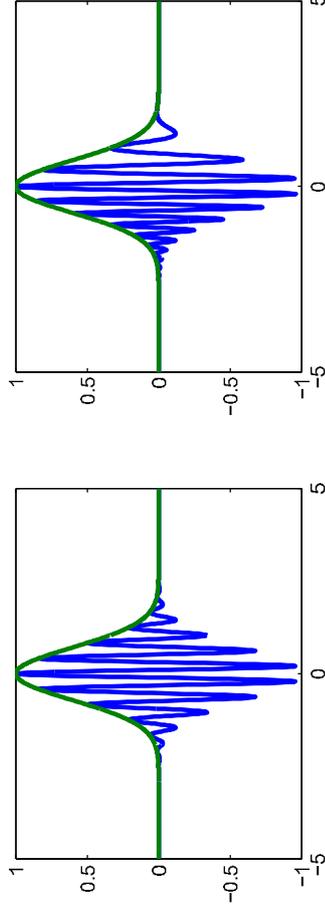
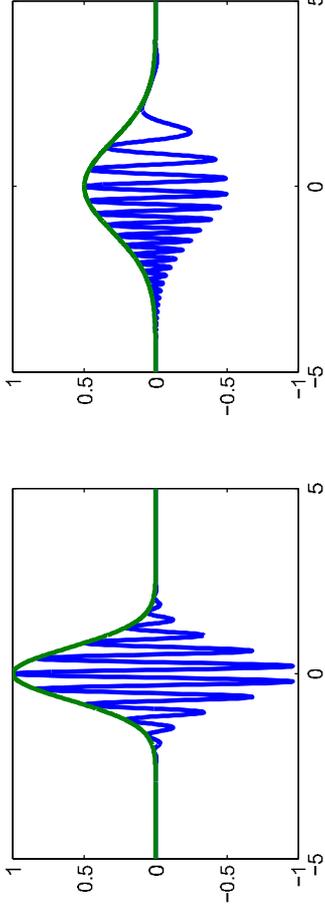


Outline

- Dispersion and nonlinearity compensation by Spectral Phase Conjugation
- Beating temporal quantum limits by Quantum Soliton Control
- Beating spatial quantum limits by Self-Focusing
- Multiphoton Absorption Rate of Quantum Lithography
- Beating resolution limits by Dielectric Slabs



Ultrashort Pulse Propagation Effects

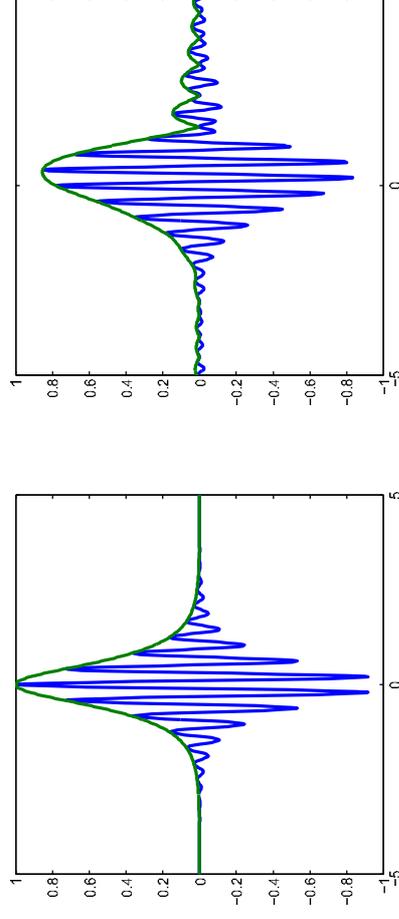


Kerr Nonlinearity ($\Delta n(t) \sim I(t)$)

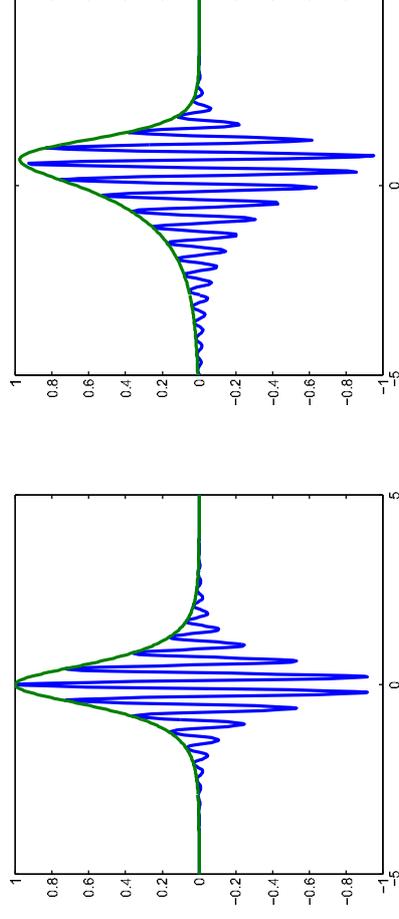


High-Order Effects

Third-Order Dispersion



Self-Steepening



Stimulated Raman Scattering, ...

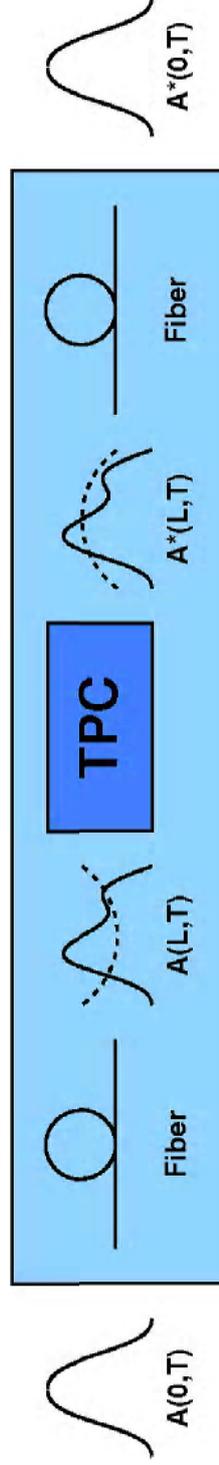
Agrawal, *Nonlinear Fiber Optics* (Academic Press, San Diego, 2001)



Ultrashort Pulse Propagation

$$\frac{\partial A}{\partial z} = -\frac{\alpha}{2}A - \frac{i\beta_2}{2}\frac{\partial^2 A}{\partial T^2} + \frac{\beta_3}{6}\frac{\partial^3 A}{\partial T^3} + i\gamma \left[|A|^2 A + \frac{i}{\omega_0} \frac{\partial}{\partial T} (|A|^2 A) - T_{RA} \frac{\partial |A|^2}{\partial T} \right] \quad (1)$$

- Perform $a(\omega) \rightarrow a^*(2\omega_0 - \omega)$, or equivalently $A(T) \rightarrow A^*(T)$, to compensate for GVD and Kerr effect and if **loss**, **higher-order dispersion**, and **self-steepening** can be neglected.
- **Temporal Phase Conjugation (TPC)**



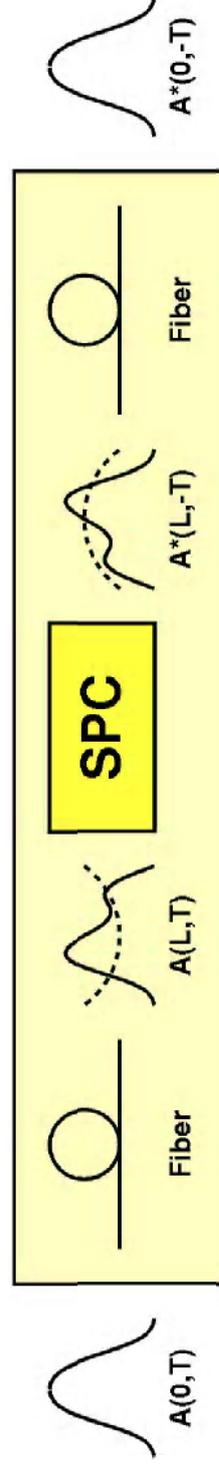
Yariv, Fekete, and Pepper, Optics Letters, **4**, 52 (1979),

Fisher, Suydam, and Yevick, Optics Letters, **8**, 611 (1983).

Spectral Phase Conjugation

$$\frac{\partial A}{\partial z} = -\frac{\alpha}{2}A - \frac{i\beta_2}{2} \frac{\partial^2 A}{\partial T^2} + \sum_{n=3}^{\infty} \frac{i\beta_n}{n!} \left(i \frac{\partial}{\partial T}\right)^n A + i\gamma \left[|A|^2 A + \frac{i}{\omega_0} \frac{\partial}{\partial T} (|A|^2 A) - T_{RA} \frac{\partial |A|^2}{\partial T} \right] \quad (2)$$

- Perform $a(\omega) \rightarrow a^*(\omega)$, or equivalently $A(T) \rightarrow A^*(T_0 - T)$, to compensate for dispersion of all orders, Kerr effect, and self-steepening if **loss** and **stimulated Raman scattering** can be neglected.
- Spectral Phase Conjugation (SPC)

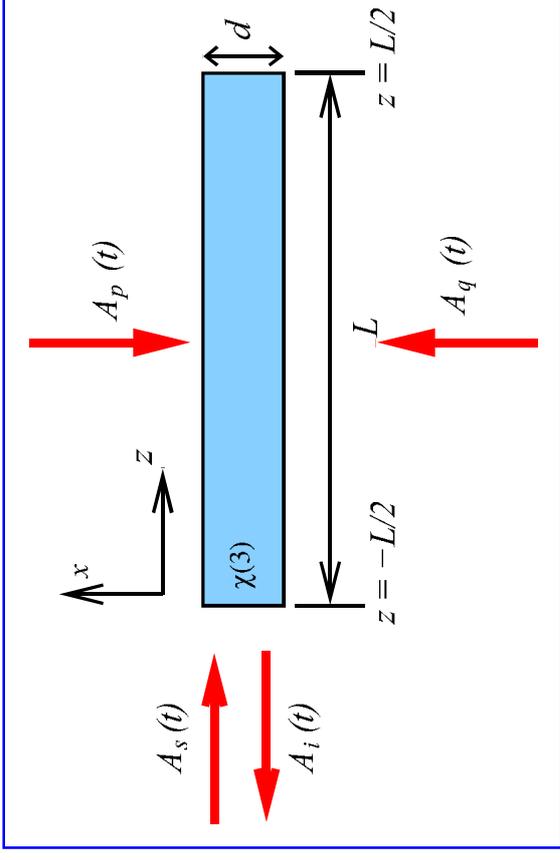


Tsang and Psaltis, Optics Letters **28**, 1558 (2003)

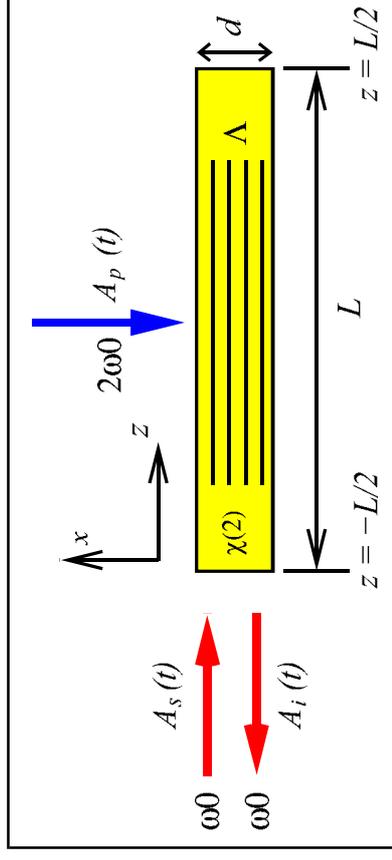


Methods of Performing SPC

● Tsang and Psaltis, Optics Express, **12**, 2207 (2004)



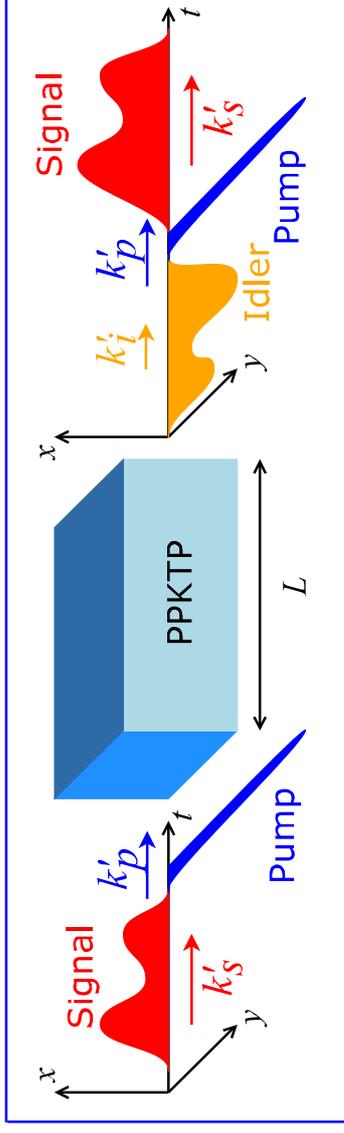
● Tsang and Psaltis, Optics Communications, **242**, 659 (2004)





SPC via Extended Phase Matching

● Tsang, JOSA B 23, 861 (2006)



- Extended Phase Matching: Quasi-Phase Matching + Group Velocity Matching
- Quasi-Phase Matching satisfied by periodic poling of nonlinear crystals
- Group Velocity Matching satisfied by material dispersion, such as KTP at 1584 nm.



Coupled-Mode Equations

• z derivatives can be neglected if the pump pulse is short enough:

$$\frac{\partial A_s}{\partial z} + k'_s \frac{\partial A_s}{\partial t} = j\chi A_p(t - k'_p z) A_i^* \quad (3)$$

$$\frac{\partial A_i^*}{\partial z} + k'_i \frac{\partial A_i^*}{\partial t} = -j\chi A_p(t - k'_p z) A_s \quad (4)$$

• Approximate solutions:

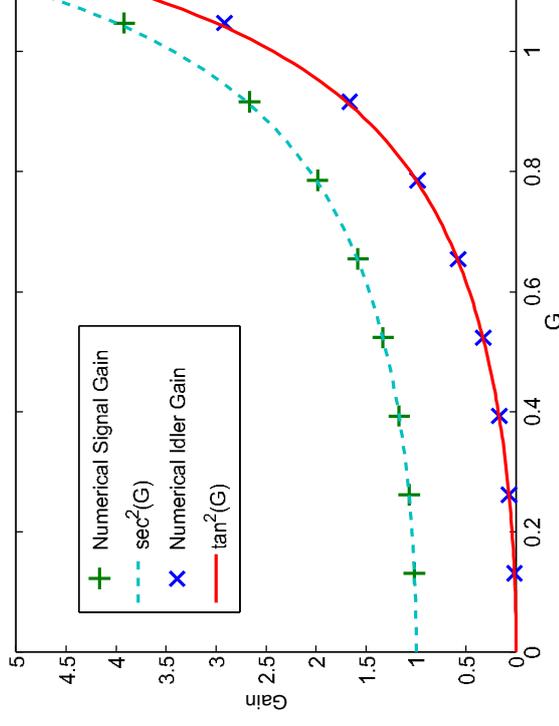
$$A_s(L, t) = A_s(0, t) \sec(G) + jA_i^*(0, -t) \tan(G) \quad (5)$$

$$A_i(L, t) = A_i(0, t) \sec(G) + jA_s^*(0, -t) \tan(G) \quad (6)$$

$$G = \frac{\chi}{\gamma} \int A_p(\tau) d\tau = \left(\frac{1}{1 - k'_p/k'_s} \right) \chi \int A_p(\tau) v d\tau \quad (7)$$



Mirrorless Parametric Oscillation

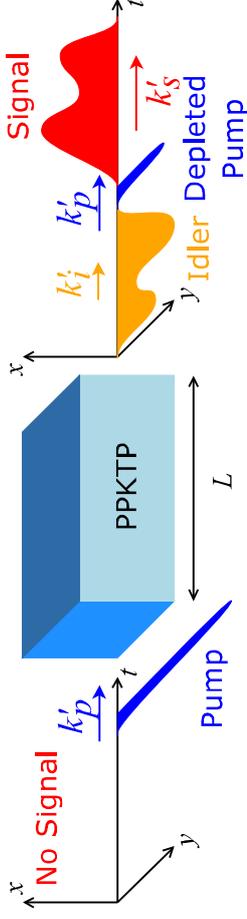


$$A_s(L, t) = A_s(0, t)\sec(G) + jA_i^*(0, -t)\tan(G) \quad (8)$$

$$A_i(L, t) = A_i(0, t)\sec(G) + jA_s^*(0, -t)\tan(G) \quad (9)$$

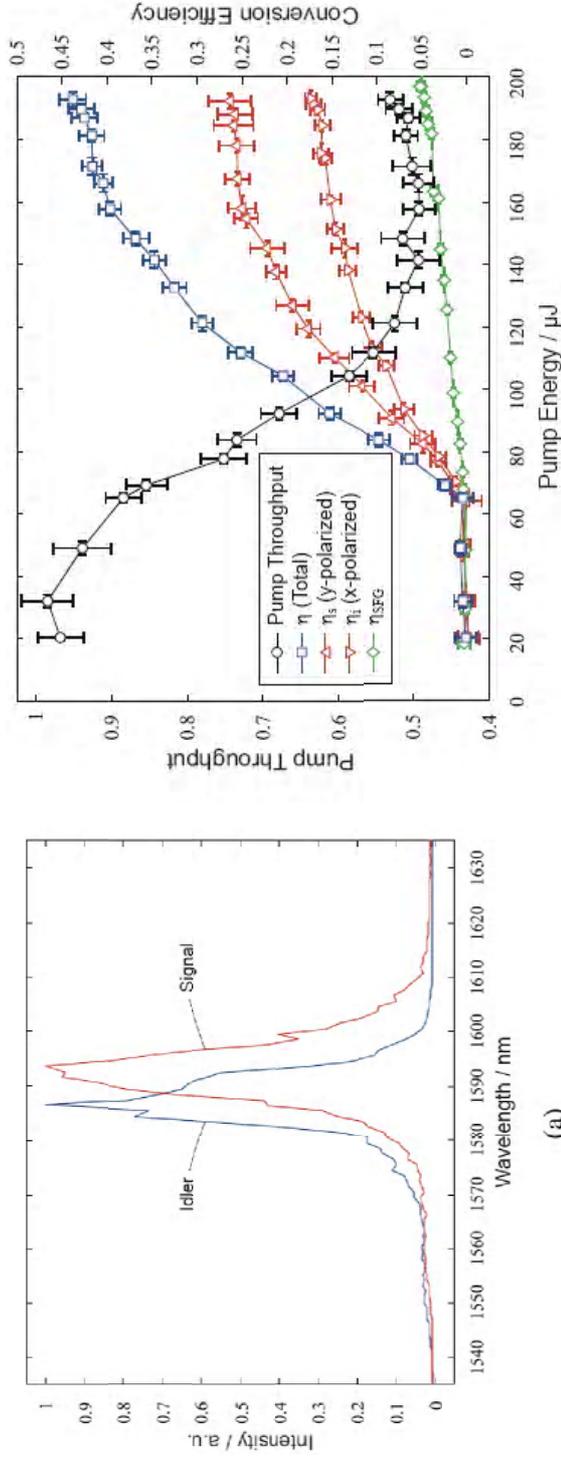
- What happens when $G = \pi/2$, and $\sec(G), \tan(G) = \infty$?
- z derivatives can no longer be neglected, gain increases exponentially with respect to z .
- Analogous to mirrorless optical parametric oscillation

Experimental Demonstration of Mirrorless OPO

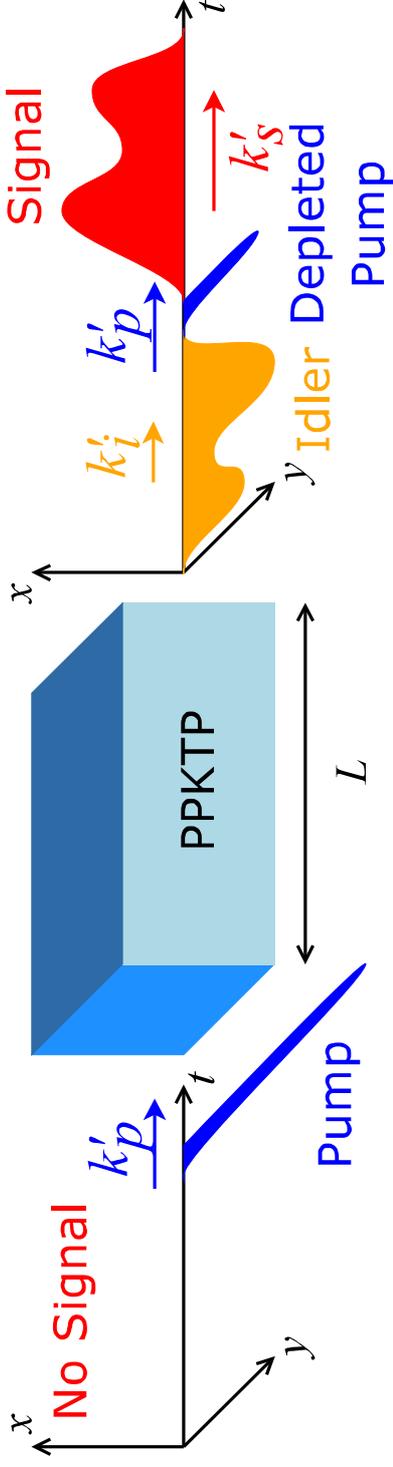


- 3cm periodically-poled KTP crystal from Raicol Crystals, dispersed femtosecond pump pulse at 792nm

- 43% down conversion efficiency, 140 dB equivalent gain



Spontaneous Parametric Down Conversion



- Classical theory predicts zero output for zero input:

$$\frac{\partial A_s}{\partial z} + k'_s \frac{\partial A_s}{\partial t} = j\chi A_p(t - k'_p z) A_i^* \quad (10)$$

$$\frac{\partial A_i^*}{\partial z} + k'_i \frac{\partial A_i^*}{\partial t} = -j\chi A_p(t - k'_p z) A_s \quad (11)$$

- Quantum theory predicts **entangled photon pair generation** even for zero input.

Giovannetti *et al.* (MIT), Physical Review Letters, **88**, 183602 (2002),

Kuzucu *et al.* (MIT), Physical Review Letters, **94**, 083601 (2005)



Comparison with Kuzucu et al.'s experiment

- Number of generated photon pairs per pump pulse is given by $\tan^2(G)$.
[Tsang, JOSAB **23**, 861 (2006)]
- Using their parameters, $\lambda_0 = 1584$ nm, $\chi^{(2)} = 7.3$ pm/V, $n_0 = 2$, $\gamma = 1.5 \times 10^{-10}$ s/m, $T_p = 100$ fs, average pump power = 350 mW, diameter = 200 μm , and pump repetition rate $f_r = 80$ MHz, the spontaneously generated photon pairs per second is theoretically given by

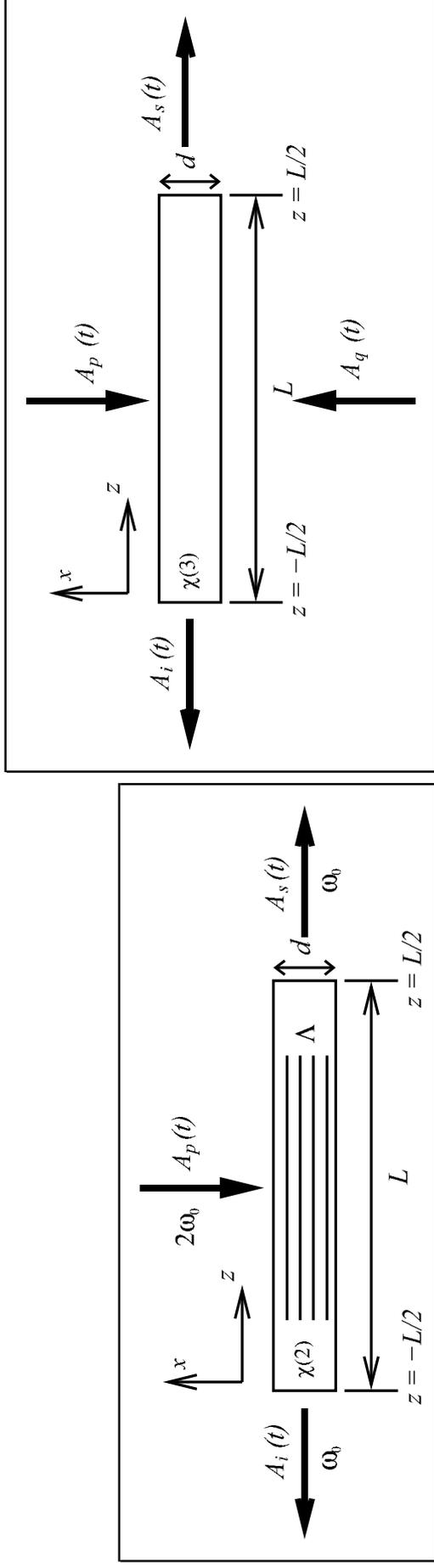
$$G = 0.2, \tag{12}$$

$$f_r \tan^2(G) = 3.6 \times 10^6/\text{s}. \tag{13}$$

Kuzucu et al. (MIT), Physical Review Letters, **94**, 083601 (2005):

coupling efficiency into the PM fiber. From the detection efficiencies and our measurement duty cycle we estimate a single spatial fiber-optic mode pair production rate of $\sim 4 \times 10^6/\text{s}$ at 350 mW of pump power.

Spontaneous Spectral Phase Conjugation



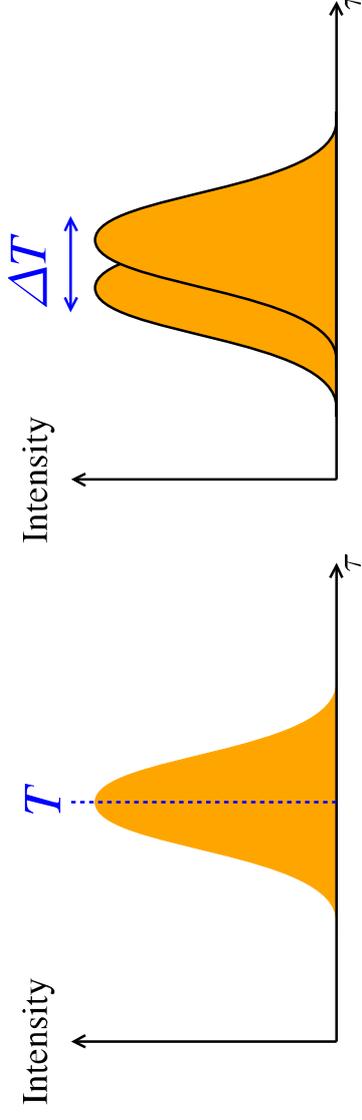
Spontaneous SPC can also create **positive frequency correlation** and **negative time correlation**.

Tsang and Psaltis, Physical Review A, **71**, 043806 (2005)



What's Special about These Photon Pairs?

- The entangled photons are **frequency correlated** and **time anti-correlated**.
- One-way autocompensating cryptography [Walton et al., PRA **67**, 062309 (2003)], **Quantum enhancement of timing accuracy** [Giovannetti et al., Nature **412**, 417 (2001)].



$$T = \frac{\int_{-\infty}^{\infty} t |A(t)|^2 dt}{\int_{-\infty}^{\infty} |A(t)|^2 dt} \quad (14)$$

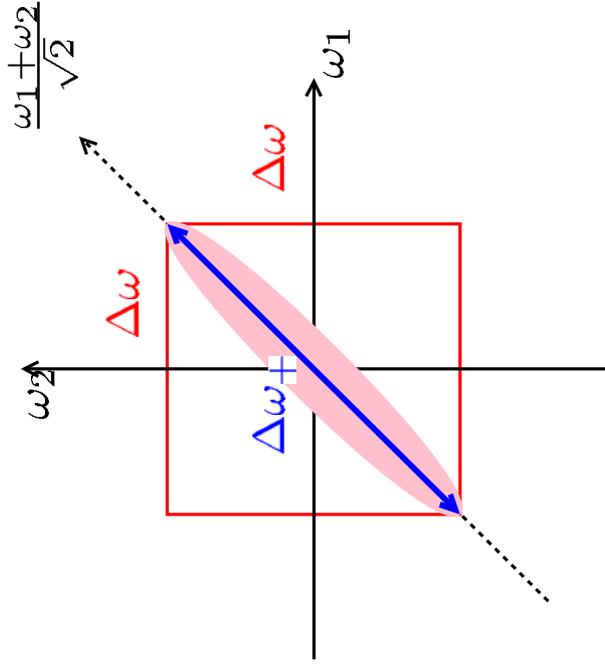
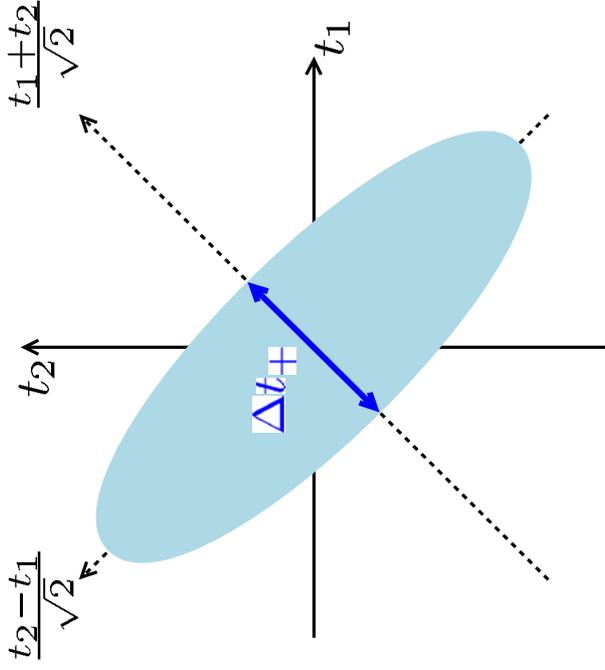
- Time-anti-correlated photons can achieve a **lower uncertainty in T for the same bandwidth**.
- Useful for **clock synchronization** or **positioning**.
- Analogous to how **mutual funds** work: selecting negatively-correlated stocks reduces risk.



Beating the SQL of Arrival Time Accuracy

Given some bandwidth limit $\Delta\omega$, (the red box), what is the smallest achievable

$$\Delta\left(\frac{t_1+t_2}{2}\right)?$$



$$\Delta\left(\frac{t_1+t_2}{2}\right) \sim \frac{1}{\sqrt{2}\Delta\omega_+} = \frac{1}{\sqrt{2}(\sqrt{2}\Delta\omega)} = \frac{1}{2\Delta\omega} \quad (15)$$

$\sqrt{2}$ enhancement with negative time correlation and positive frequency correlation.



Multiphoton Enhancement

Giovannetti, Lloyd, and Maccone, Nature **412**, 417 (2001)

- N independent photons:

$$\Delta T \geq \frac{1}{2\sqrt{N}\Delta\omega} \quad (\text{Standard Quantum Limit}) \quad (16)$$

- e.g.: $W = 100 \text{ ps}$, $N = 10^{10}$, $\Delta T = 1 \text{ fs}$

- N negatively-time-correlated photons:

$$\Delta T \geq \frac{1}{2N\Delta\omega} \quad (\text{Ultimate Quantum Limit}) \quad (17)$$

- $N = 2$ is quite useless compared to $N \gg 1$

- How to create multiphoton time anti-correlation with $N \gg 1$?

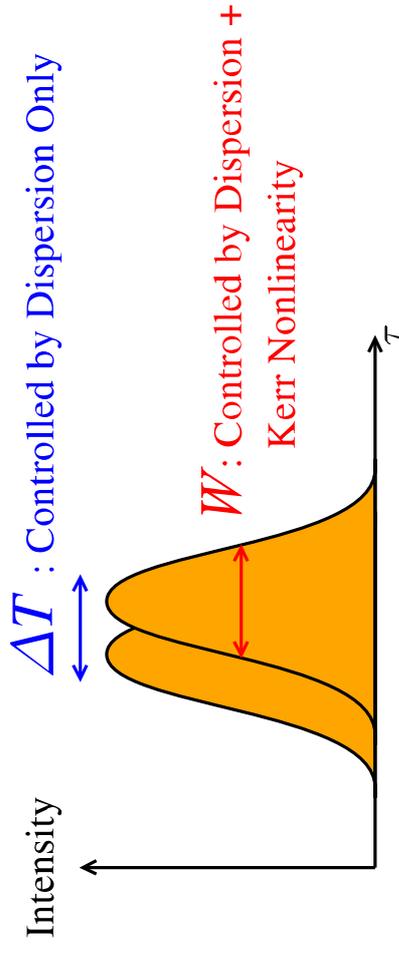


Quantum Theory of Optical Fiber Solitons

- **Classical theory:** soliton is a stable solitary wave due to balance between **anomalous GVD** and **Kerr effect**



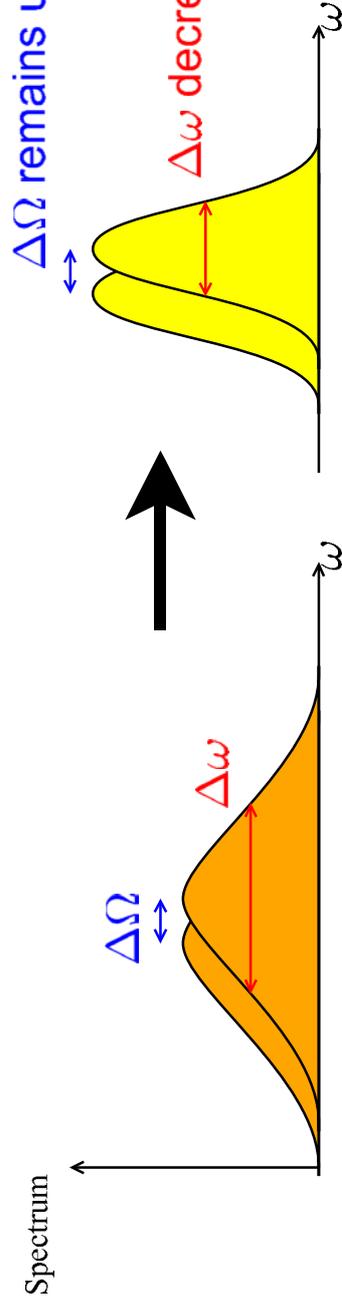
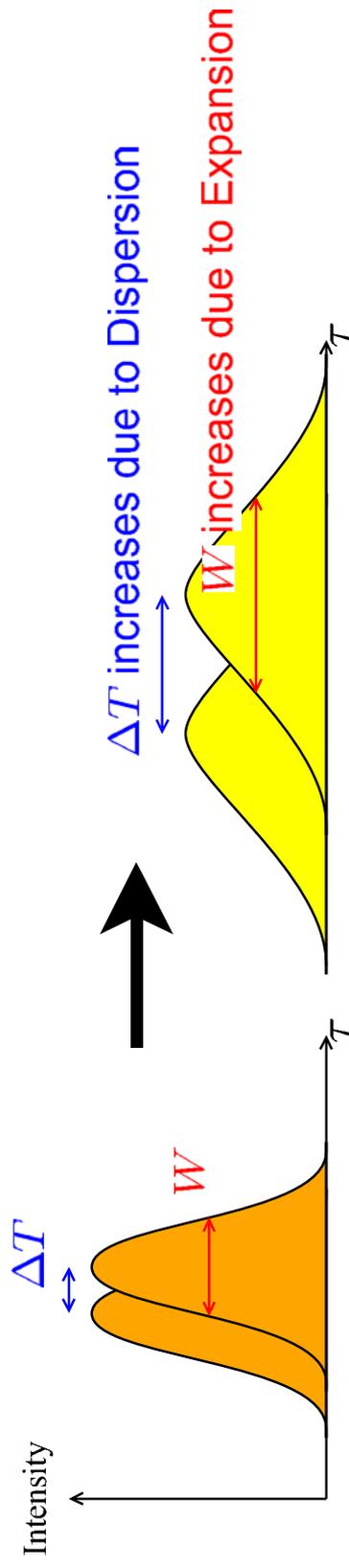
- **Quantum theory:** Stable pulse shape and bandwidth due to balance between GVD and Kerr effect, but the **average position of the pulse is affected by dispersion only.**





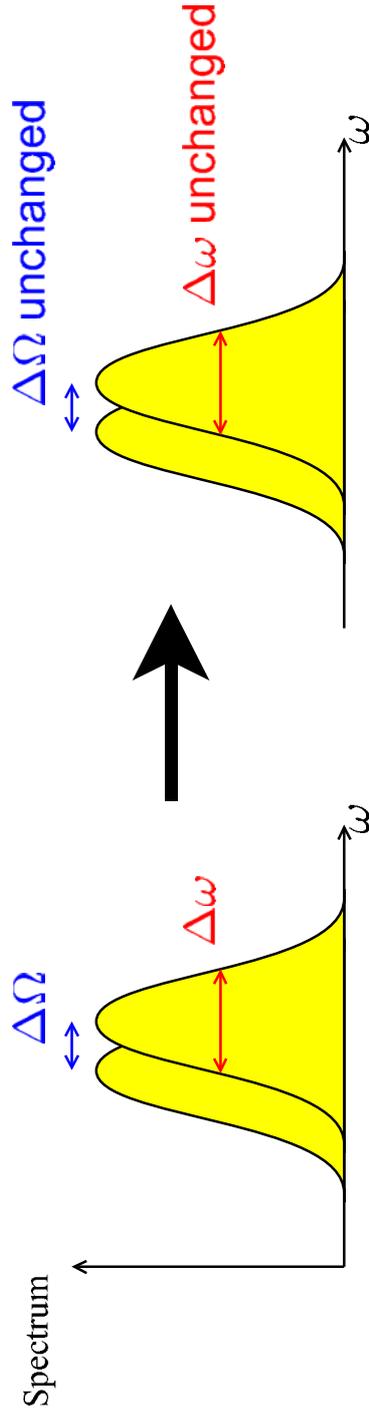
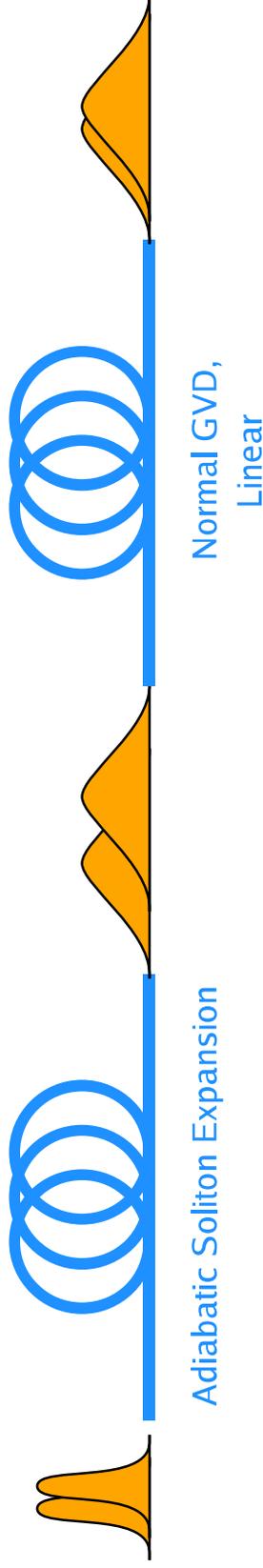
Adiabatic Soliton Expansion

1. Adiabatically reduce the **Kerr nonlinearity** or increase the **group-velocity dispersion** along the fiber



Quantum Dispersion Compensation

- 2. Compensate for dispersion of T in a second fiber with $b'L' = -bL$

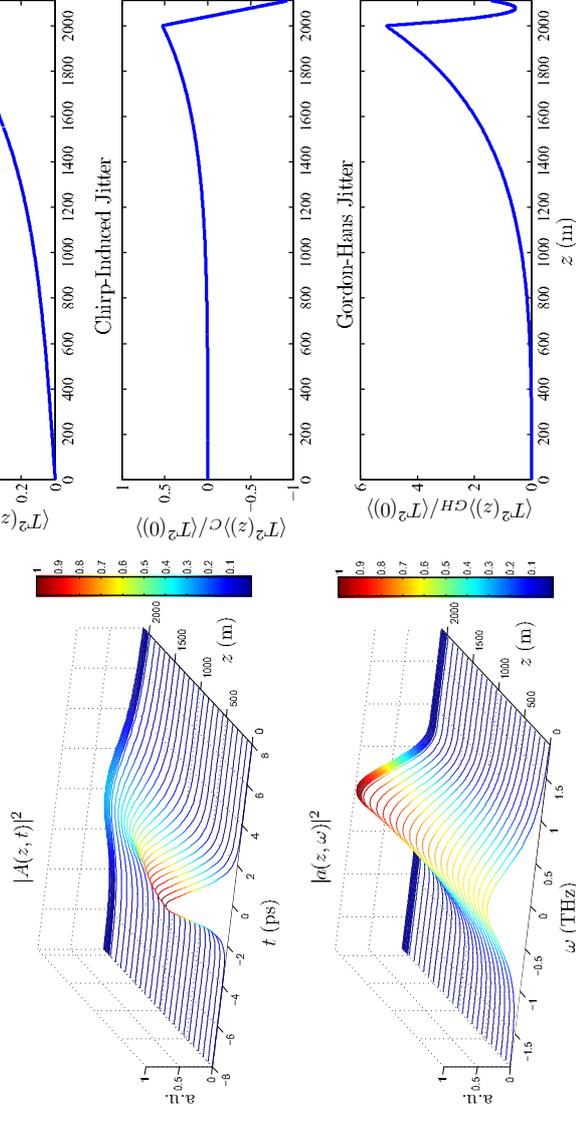


- ΔT is the same as the input, but $\Delta\omega$ is reduced, so $\Delta T < 1/(2\sqrt{N}\Delta\omega)$.
- Subfemtosecond timing jitter detection can be performed by cross-correlation measurements via sum-frequency generation or balanced homodyne measurements with a reference local oscillator pulse.



Decoherence of Adiabatic Soliton Control

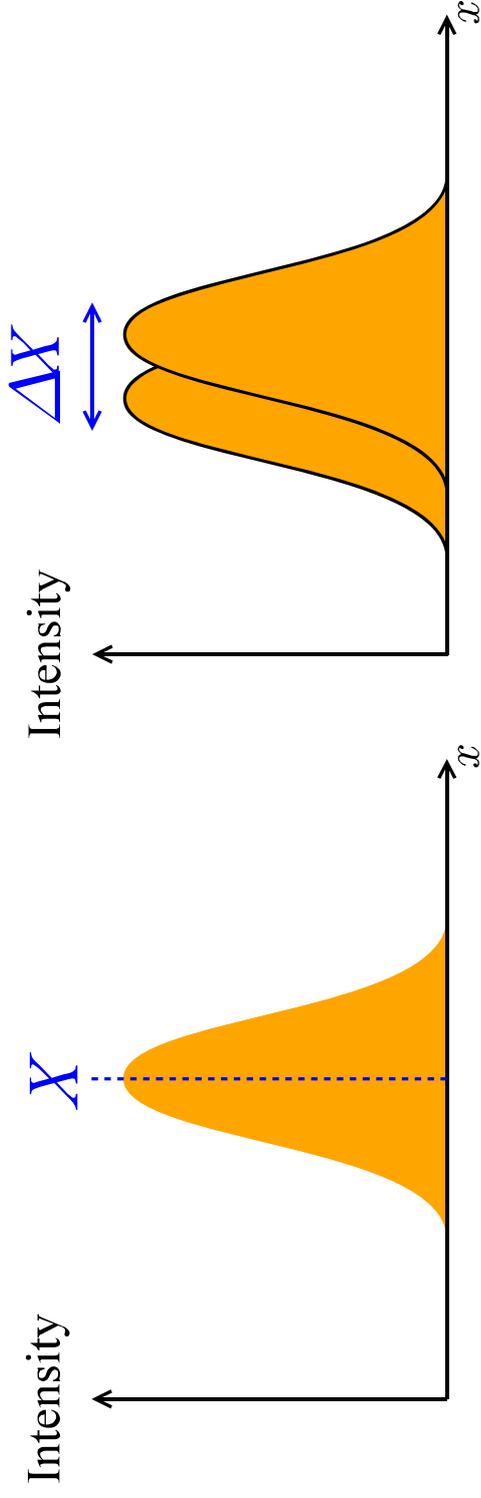
- Decoherence mainly due to **loss of photons**. Main source of quantum noise is the **Gordon-Haus timing jitter**.
- The decoherence effect can be estimated numerically using a semiclassical analysis:



- Using realistic parameters, the achievable squeezing is predicted to be ~ 3.8 dB.
- Not great, but enough for proof-of-concept demo.



Optical Beam Position



$$\Delta X_{SQL} \sim \frac{1}{2\sqrt{N}\Delta k}, \quad \Delta X_{UQL} \sim \frac{1}{2N\Delta k} \quad (18)$$

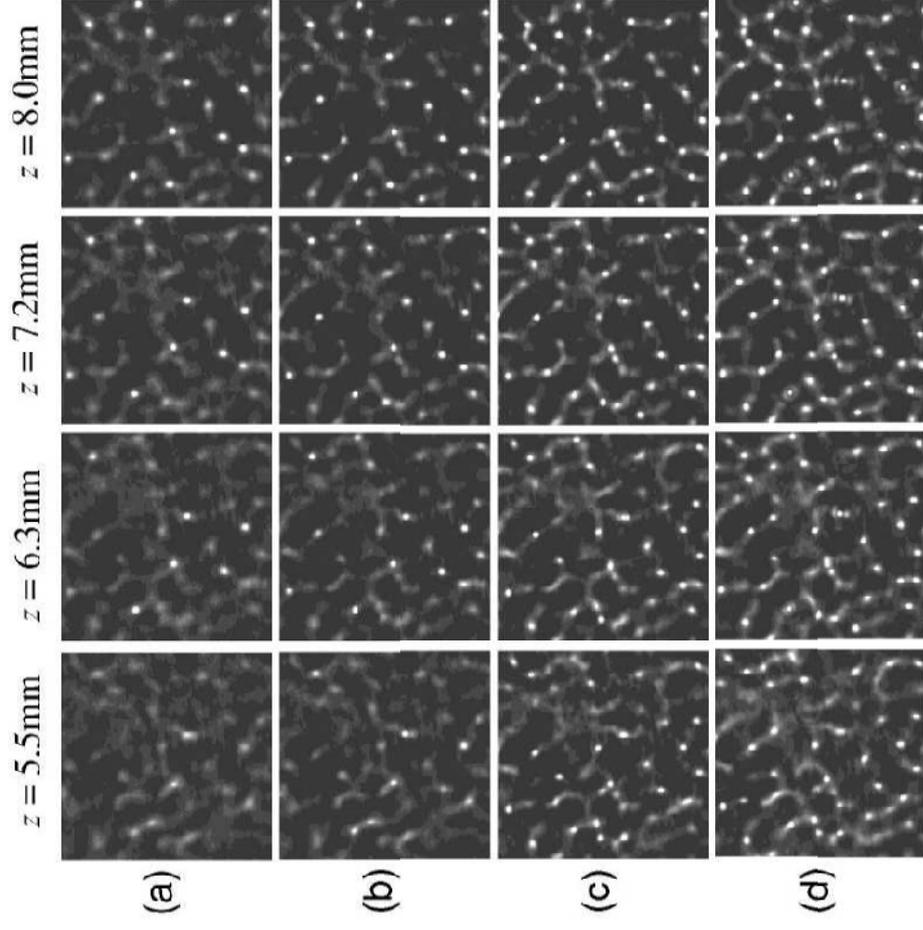
Important for

- Nanoparticle Detection
- Atomic Force Microscope
- Lithography



2D Self-Focusing Collapse

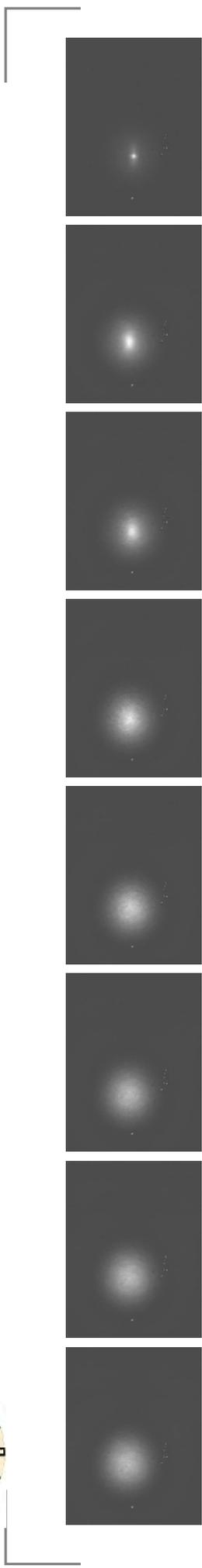
Balance between diffraction and Kerr effect is **unstable**.



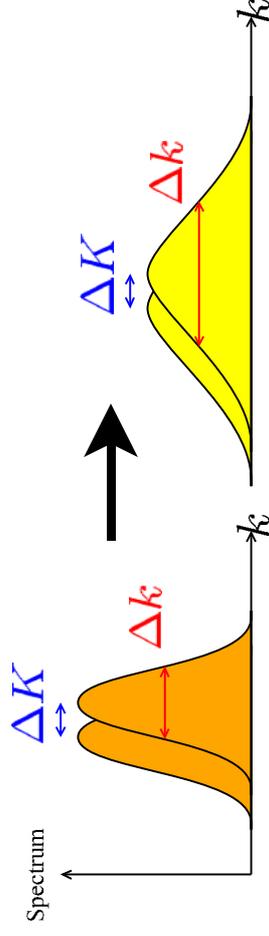
Centurion, Pu, Tsang, and Psaltis, *Physical Review A* **71**, 063811 (2005)



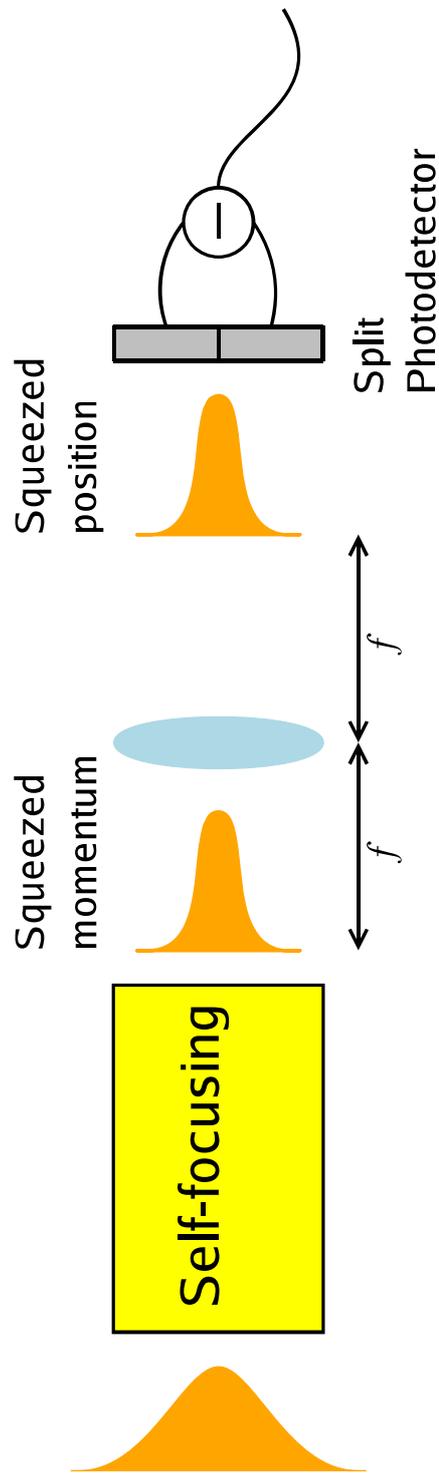
Spatial Quantum Enhancement by Self-Focusing



increasing pulse energy \rightarrow



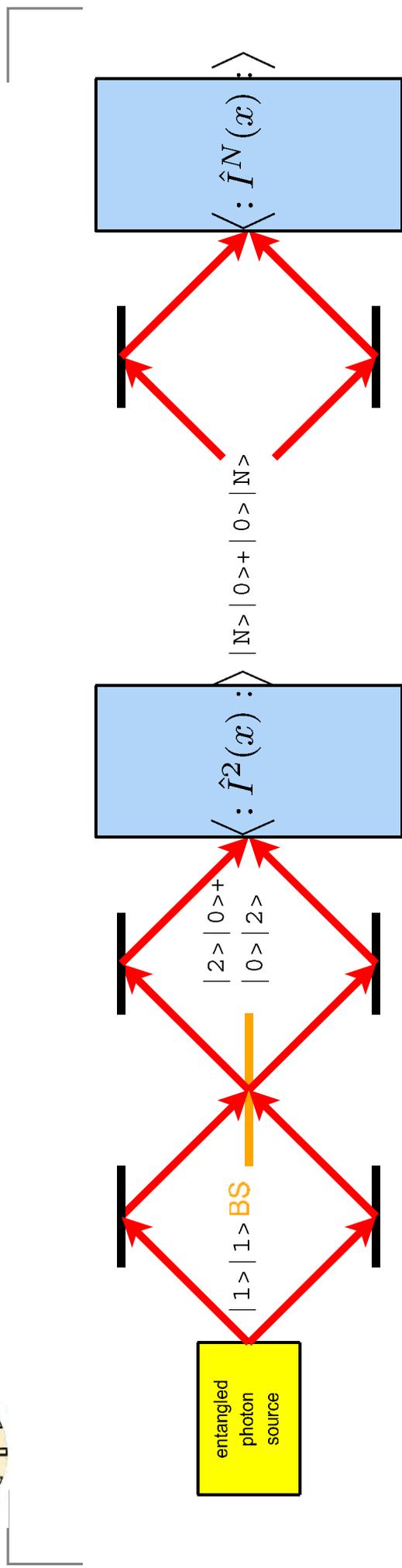
Use a **Fourier-transform lens** to transform to real space



e.g. $W = 3 \text{ mm}$, $N = 10^{10}$, $\Delta X = 30 \text{ nm}$



Quantum Lithography



$$\langle : \hat{I}^2(x) : \rangle \sim \cos^2(2k_x x), \quad \langle : \hat{I}^N(x) : \rangle_{NOON} \sim \cos^2(N k_x x), \quad (19)$$

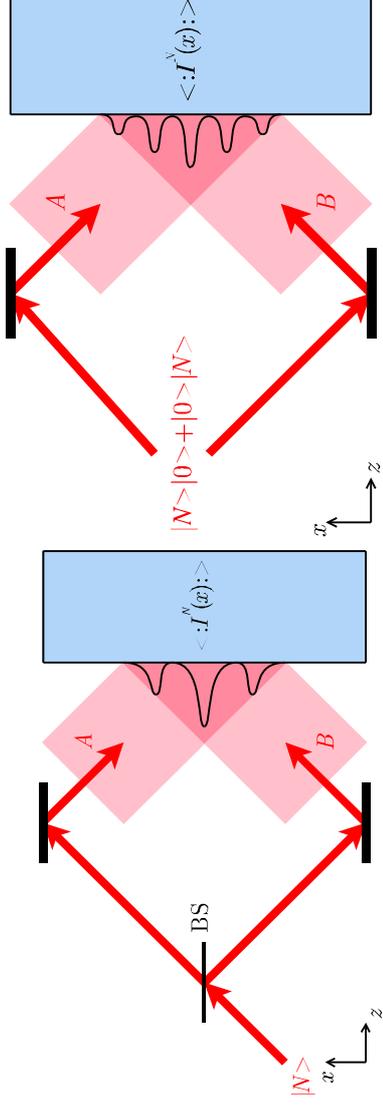
Factor-of- N resolution enhancement compared to one-photon absorption or classical multiphoton lithography:

$$\langle : \hat{I}(x) : \rangle \sim \cos^2(k_x x), \quad \langle : \hat{I}^N(x) : \rangle_{classical} \sim \cos^{2N}(k_x x) \quad (20)$$

Boto et al., Physical Review Letters 85, 2733 (2000).



Multiphoton Absorption Rate of NOON State



$$\langle : \hat{I}^N(x) : \rangle_{classical} \sim \cos^{2N}(\kappa x) \quad \langle : \hat{I}^N(x) : \rangle_{NOON} \sim \cos^2(N\kappa x) \quad (21)$$

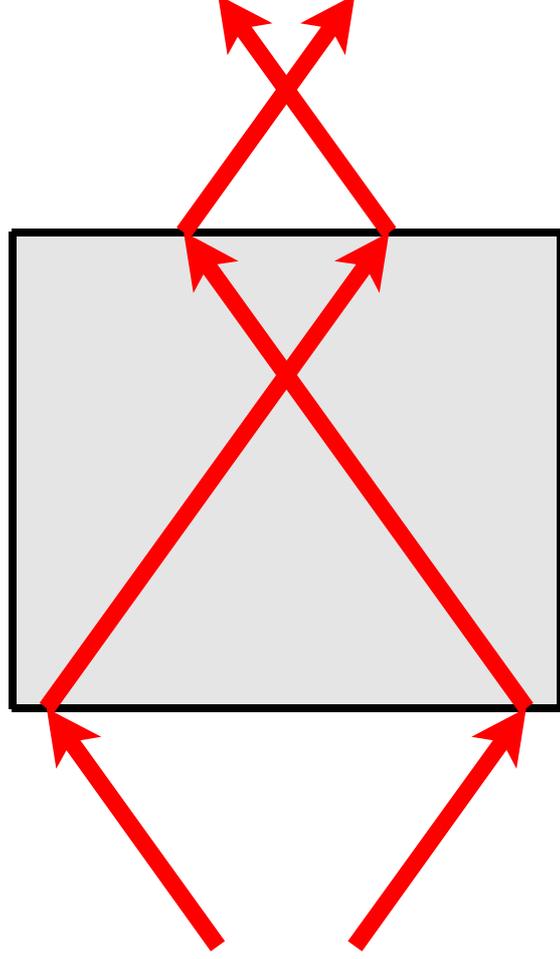
$$\langle : \hat{I}^N(0) : \rangle_{classical} = 2^{N-1} \langle : \hat{I}^N(0) : \rangle_{NOON} \quad (22)$$

- Spatial frequency correlation leads to **position anti-correlation**
- Photons arrive further apart from one another, less chance of all photons arriving at the same atom.
- Proof-of-concept demo of quantum lithography reported in D'Angelo *et al.*, PRL **87**, 013602 (2001), so the reduced coincidence rate can be readily tested experimentally.

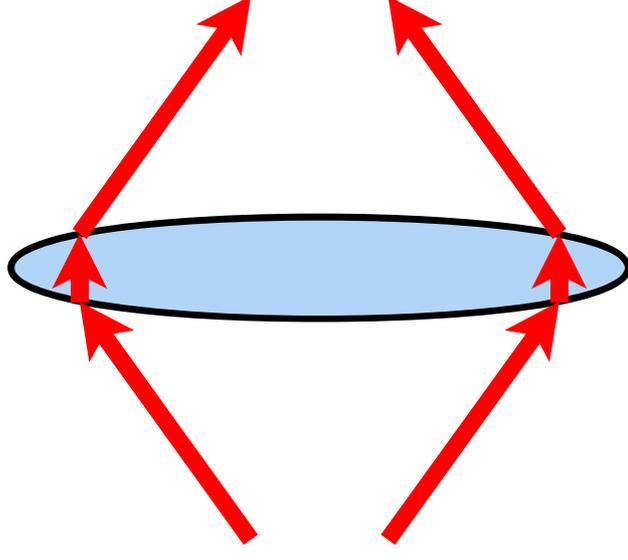


Veselago: Negative Refraction

$$n = -1$$



$$n > 1$$

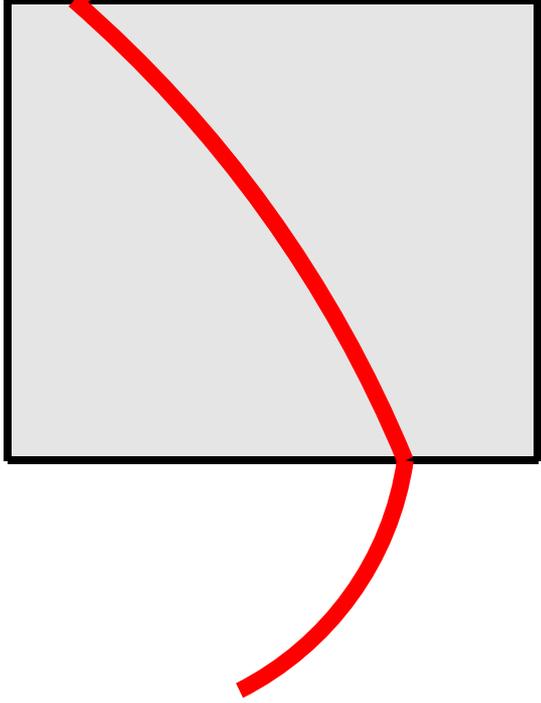


Veselago, Sov. Phys. Usp. **10**, 509 (1968)

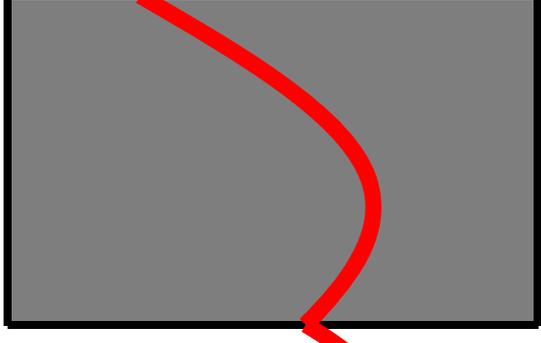


Pendry: Evanescent Wave Amplification

$$n = -1$$



$$\epsilon < 0$$



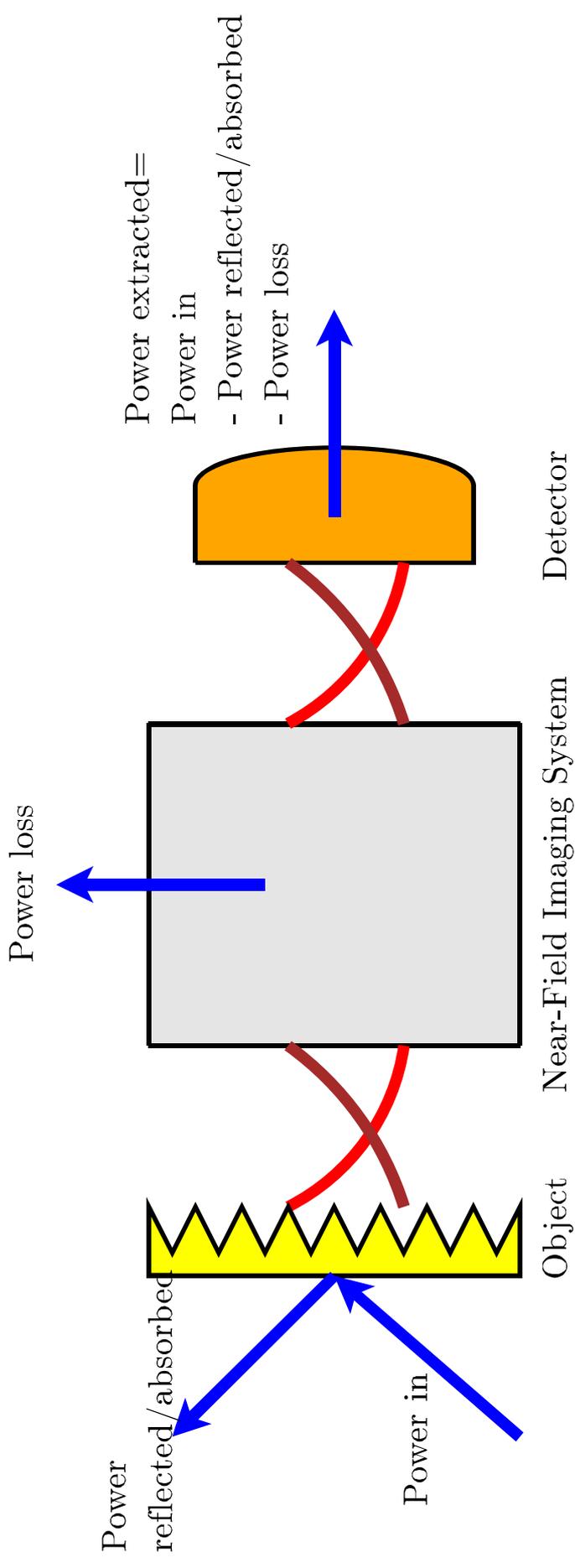
$$T \rightarrow \exp(-ik_z d), \quad R \rightarrow 0$$

(23)

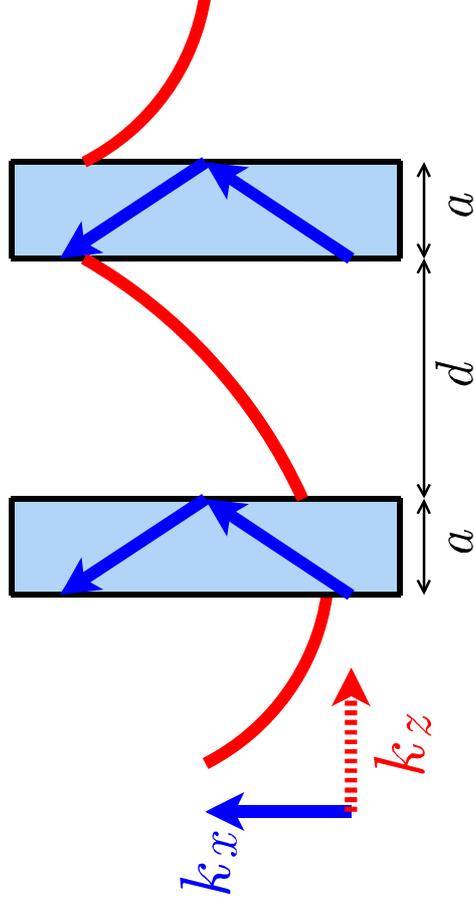
Pendry, Phys. Rev. Lett. **85**, 3966 (2000)



Importance of Low Loss



Two Dielectric Slabs



$$R = \Gamma + \frac{\tau^2 \Gamma \exp(2ik_z d)}{1 - \Gamma^2 \exp(2ik_z d)} = 0, \quad T = \frac{\tau^2 \exp(ik_z d)}{1 - \Gamma^2 \exp(2ik_z d)} = -\exp(-ik_z d) \quad (24)$$

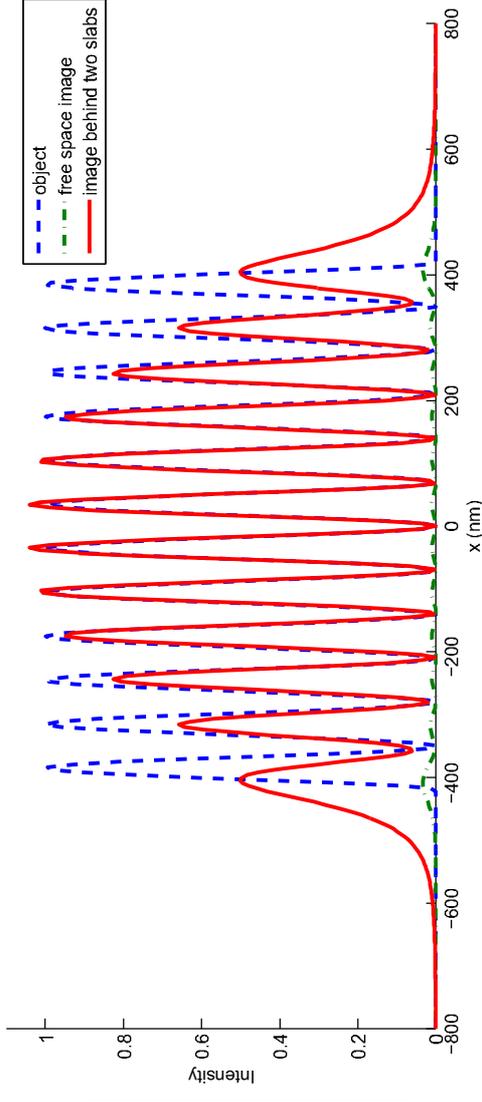
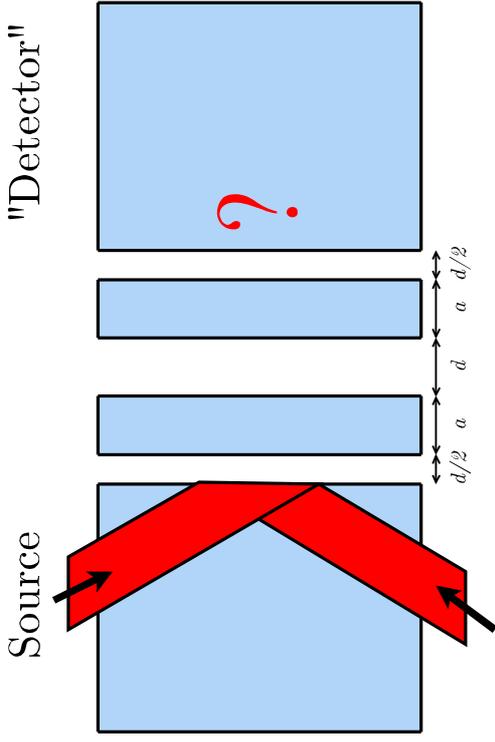
for some k_x .

Tsang and Psaltis, Optics Letters, **31**, 2741 (2006); Erratum: **32**, 86 (2006).



Numerical Example

• $\lambda = 230 \text{ nm}$, $n = 2.7$, $a = 20 \text{ nm}$, $d = 20 \text{ nm}$, TE polarization,



- Low loss
- Many spatial modes
- High refractive index material available (transparent down to $\lambda = 230 \text{ nm}$, $n = 2.7$ for diamond)
- non-contact imaging, suitable for lithography and bio-imaging



Miscellaneous

- Correspondence between Nonlinear Optics and Fluid Dynamics
[Tsang and Psaltis, e-print physics/0604149](#)
- Reverse propagation of femtosecond pulses in optical fiber (collaboration with
Fiorenzo Omenetto at Tufts)
[Tsang, Psaltis, and Omenetto, Optics Letters 28, 1873 \(2003\)](#)
- Filamentation in Kerr Medium
[Centurion, Pu, Tsang, and Psaltis, Physical Review A 71, 063811 \(2005\)](#)
- Quantum Temporal Imaging
[Tsang and Psaltis, Physical Review A 73, 013822 \(2006\)](#)

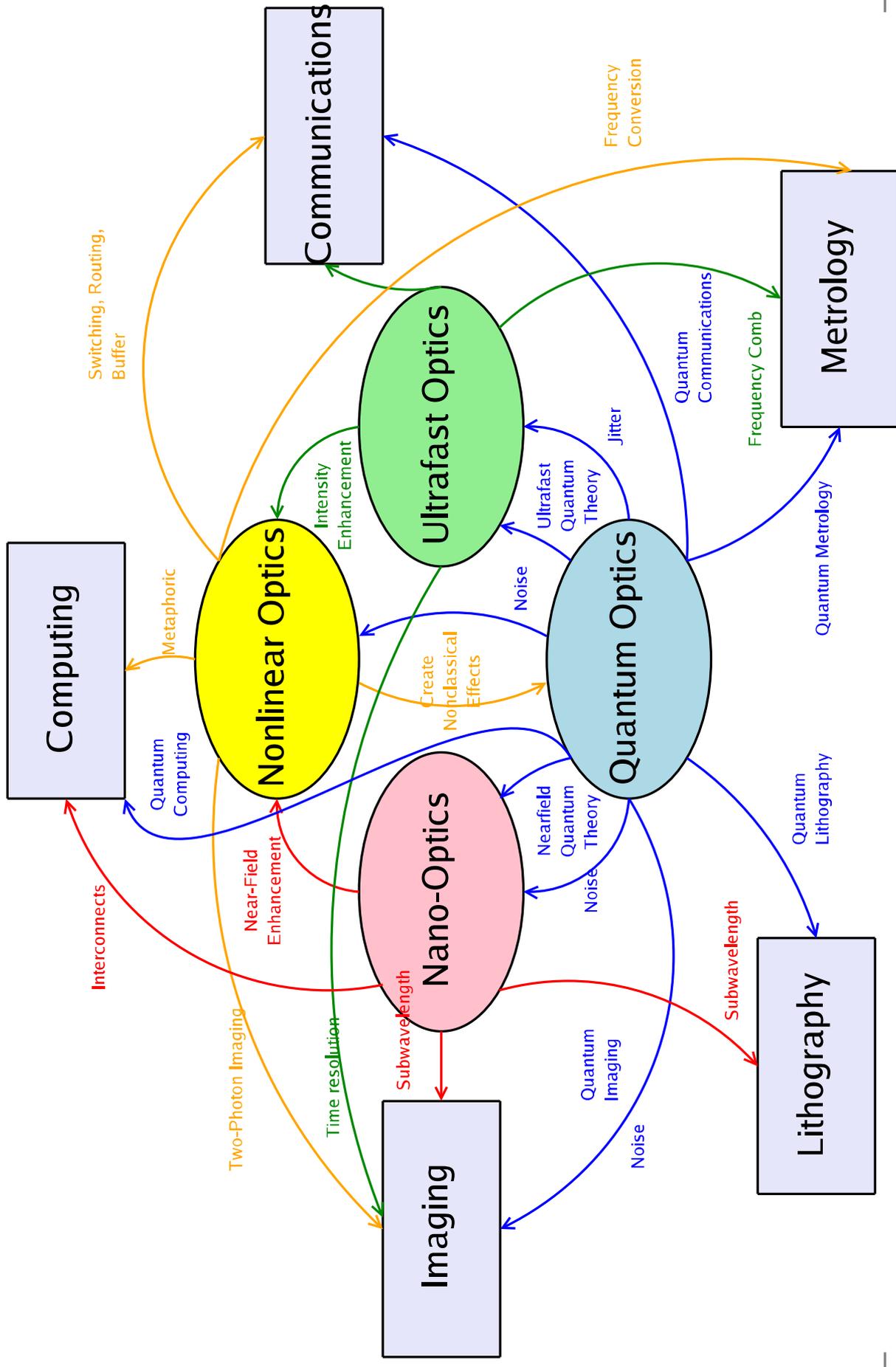


Future Work

- Quantum theory of mirrorless optical parametric oscillators
- Quantum information processing via scalar and vector solitons
- Spatial quantum information processing via spatial solitons
- Effect of loss and decoherence
- Quantum limits on spatial, temporal, and spectral information capacity of optical fields
- Beating the resolution limit of λ/n by the use of dielectrics, e.g. photonic crystals, coupled resonators
- Quantum near-field optics
- Correspondence between nonlinear optics and viscous fluid dynamics
- Application to Bose-Einstein condensates and superfluids
- Experiments



Quantum Optical Engineering





Publications

1. **M. Tsang**, "Relationship between resolution enhancement and multiphoton absorption rate in quantum lithography," *Physical Review A* **75**, 043813 (2007).
2. **M. Tsang** and **D. Psaltis**, "Reflectionless evanescent-wave amplification by two dielectric planar waveguides," *Optics Letters* **31**, 2741 (2006) [Erratum: **32**, 86 (2006)].
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1. **Y. Pu**, **J. Wu**, **M. Tsang**, and **Psaltis**, "Ultrafast mirrorless optical parametric oscillator in periodically poled KTiOPO4 via extended phase matching," CLEO/QELS, May 2007, paper CMB5.
2. **M. Tsang** and **D. Psaltis**, "Reflectionless evanescent wave amplification by two dielectric slabs," OSA Annual Meeting, Oct 2006, paper FMB4.
3. **M. Tsang** and **D. Psaltis**, "Quantum lithography has a reduced multiphoton absorption rate," OSA Annual Meeting, Oct 2006, paper LWH4.
4. **M. Tsang** and **D. Psaltis**, "Quantum temporal imaging," CLEO/QELS, May 2006, paper QWB5.
5. **M. Centurion**, **M. Tsang**, and **D. Psaltis**, "Nonlinear signal processing," Invited Talk, LEOS Annual Meeting, Oct 2005, paper TuEE2.
6. **M. Tsang** and **D. Psaltis**, "Metaphoric optical computing of fluid dynamics," CLEO/QELS, May 2005, paper QML6.
7. **M. Centurion**, **Y. Pu**, **M. Tsang**, and **D. Psaltis**, "Phase transition in the filament generation process in a Kerr medium," CLEO/QELS, May 2005, paper QMI3.
8. **M. Tsang** and **D. Psaltis**, "Metaphoric optical computing for fluid dynamics," Invited Paper, Proceedings of SPIE, 5735, 1 (Apr 2005).
9. **M. Tsang** and **D. Psaltis**, "Spectral phase conjugation with cross-phase modulation compensation," OSA Annual Meeting, Oct 2004, paper FWH44.

Preprints:

1. **M. Tsang**, "Decoherence of quantum-enhanced timing accuracy," e-print arXiv:0704.0663 [submitted].
2. **M. Tsang** and **D. Psaltis**, "Metaphoric optical computing of fluid dynamics," e-print physics/0604149.

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